Random organization with mediated interactions

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Reversible-irreversible transitions (RIT)

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Stroboscopic time: continuous phase transition





Reversible-irreversible transitions (RIT)

Dense suspensions [Pine et al, Corté et al, 2005-2009]

Emulsions, moderately concentrated [Jeanneret & Bartolo, 2014-2015], very concentrated [Knowlton, Pine & Cipelletti, 2014]

Soft glasses, experimental [Hima Nagamanasa et al, 2014], simulations [Fiocco, Foffi & Sastry, 2013]

Dry granular material [Royer & Chaikin, 2015]





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Stroboscopic sampling:

absorbing phase transition with infinitely many absorbing states and a conserved quantity

Expectation: continuous transition in the **Conserved Directed Percolation (CDP) class** [Pastor-Satorras & Vespignani, 2000]

Minimal model: random organization

[Corté, Chaikin, Gollub, Pine, Nature 2008] [Hexner & Levine, PRL 2015] [Tjhung & Berthier, PRL 2015]

Discrete-time model for stroboscopic dynamics



Random jumps with size δ Particle volume fraction $\phi \leftrightarrow$ strain amplitude γ_0

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Discrete-time model for stroboscopic dynamics



Random jumps with size δ Particle volume fraction $\phi \leftrightarrow$ strain amplitude γ_0 Absorbing phase transition $A \sim (\phi - \phi_c)^{\beta}$

 β < 1, values compatible with CDP (although debated)



Conserved directed percolation



[Menon & Ramaswamy, 2009]

Two fields:

- Local volume fraction ϕ (conserved)
- Local active particle density A (not conserved)

 $\begin{aligned} \partial_t \phi &= D_\rho \nabla^2 A\\ \partial_t A &= (-\alpha + k\phi)A - \lambda A^2 + D_A \nabla^2 A + \sigma \sqrt{A}\eta \end{aligned}$

Activity drives the dynamics

Activity creates and destroys activity, and acts like a local temperature

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Concave transtion $\langle A \rangle \sim (\phi - \phi_c)^{\beta}$ with $\beta \approx 0.64$ for d = 2 (1 in MF)

Fluctuations diverge $N\langle (A - \langle A \rangle)^2 \rangle \sim (\phi - \phi_c)^{-\gamma'}$ with $\gamma' \approx 0.37$ for d = 2 (0 in MF)

RIT is conserved directed percolation?

CDP class: density of active particles $A \sim (\gamma_0 - \gamma_c)^{\beta}$, with $\beta \approx 0.84$ for d = 3

Early experiments argued compatible with CDP [Corté et al, Nat. Phys. 2009; Hima Nagamanasa et al., PRE 2014]



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However many **non-CDP behaviors** observed in other experiments, e.g. emulsions with $\beta > 1$ (convex transition)

Semi-dilute [Weijs, Jeanneret, Dreyfus, Bartolo, PRL 2015]



Concentrated [Knowlton, Pine, Cipelletti, Soft Matter 2014]



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 Δ/Δ^{\star}

2

0.6

0.4

0.2

0



Concentrated [Knowlton, Pine,

Not clarified by simulations of realistic models either (observed CDP, continuous but not CDP, first order)





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Passive particles can spontaneously create activity (even far from other active zones)

$$\partial_t A = (-\alpha + k\phi)A - \lambda A^2 + ?? + D_A \nabla^2 A + \sigma \sqrt{A}\eta$$

Numerical simulations



Convex transition, $\beta > 1$

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Vanishing fluctuations, $\gamma' < 0$

Numerical simulations





Convex transition, $\beta > 1$

Vanishing fluctuations, $\gamma' < 0$

Activity created from passive particles motion $\partial_t A = (-\alpha + k\phi)A - \lambda A^2 + \alpha_p \phi A - \mu \phi A^{3/2} + D_A \nabla^2 A + \sigma \sqrt{A}\eta$

Mean-field is convex, with $\beta = 2!$

New universality class for RIT with mediated interactions?

Vanishing but large enough to get non-mean field behavior, Ginzburg criterion violated: $N\langle\delta A^2\rangle \gg \langle A\rangle$ close to ϕ_c

Large β linked to small fluctuations via hyperscaling $\gamma' = d\nu_{\perp} - 2\beta$ [Lübeck 2004]

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Another convex absorbing phase transition? Experts in this audience...

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Yielding of yield-stress fluids!

Hershel-Bulkley law: $\sigma = \sigma_{\rm c} + K_{\gamma}^{\cdot 1/\beta}$

Control parameter σ , order parameter $\dot{\gamma}$: $\dot{\gamma} \sim (\sigma - \sigma_c)^{\beta}$, with $\beta \in [1.5, 2]$

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Work with Tristan Jocteur, Kirsten Martens, Shana Figueiredo For Picard model

[Picard, Ajdari, Lequeux & Bocquet, 2004]

